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# A SIMPLE WAY TO ESTABLISH THE EQUATION OF SHELLS's YIELD SURFACE

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**ABSTRACT:** A yield surface is a five-dimensional surface in the six-dimensional space of stresses. The yield surface is usually convex and the state of stress of inside the yield surface is elastic. When the stress state lies on the surface the material is said to have reached its yield point and the material is said to have become plastic. Further deformation of the material causes the stress state to remain on the yield surface, even though the surface itself may change shape and size. This is because stress states that lie outside the yield surface are non-permissible in rate-independent plasticity, though not in some models of viscoplasticity. This paper deals with the yield surface of shells.

**KEYWORDS:** yield surface, shells, elastic, plastic, viscoplasticity

## INTRODUCTION

Yield surface is usually expressed and visualized in three-dimensional space of principal stresses  $(\sigma_1, \sigma_2, \sigma_3)$ , a two or three dimensional space extended by the constant stress  $(I_1, J_2, J_3)$  or a version of three-dimensional space that is Haigh-Westergaard stress. Yield surface is a frequent presence on Tresca yield surface, Huber-Von Mises yield surface, Mohr-Coulomb yield surface, Drucker-Prager yield surface, Bresler-Pister yield surface, Warnke-William yield surface, ... (Figure H1, H2 & H3)



#### **RELATIONSHIP BETWEEN STRESS, DEFORMATION AND DISPLACEMENT**

The internal forces in the shell that known as:  $N_{\theta}$ ,  $N_{\omega}$ ,  $M_{\theta}$ ,  $M_{\varphi}$  and Q.

Assumptions: (i) cross-section remains plane during deformation, (ii) the impact of negligible shear is ignored.

$$N_{\theta} = \int_{-\frac{h}{2}}^{+\frac{h}{2}} \sigma_{\theta} dz; N_{\varphi} = \int_{-\frac{h}{2}}^{+\frac{h}{2}} \sigma_{\varphi} dz; M_{\theta} = \int_{-\frac{h}{2}}^{+\frac{h}{2}} \sigma_{\theta} z dz; M_{\varphi} = \int_{-\frac{h}{2}}^{+\frac{h}{2}} \sigma_{\varphi} z dz$$

$$\varepsilon_{\theta} = \lambda_{\theta} + z\chi_{\theta}; \lambda_{\theta} = \frac{1}{r_{2}} (\operatorname{vcot} g\varphi - w); \chi_{\theta} = \frac{-\operatorname{cot} g\varphi}{r_{1}r_{2}} (v + w')$$

$$\varepsilon_{\varphi} = \lambda_{\varphi} + z\chi_{\varphi}; \lambda_{\varphi} = \frac{1}{r_{1}} (v' - w); \chi_{\varphi} = -\frac{1}{r_{1}} \left[ \frac{1}{r_{1}} (v + w') \right]$$

$$\begin{cases} \left( r_{0}N_{\varphi} \right)' - r_{1}N_{\theta} \cos\varphi - r_{0}Q + r_{0}r_{1}P_{\varphi} = o \\ \left( r_{0}N_{\varphi} \right)' - r_{1}M_{\theta} \cos\varphi - r_{0}r_{1}Q = o \end{cases}$$

Equations:

### YIELD CONDITION

Huber-Mises's condition:

$$(\sigma_1 - \sigma_2)^2 + (\sigma_2 - \sigma_3)^2 + (\sigma_3 - \sigma_1)^2 = 2\sigma_0^2$$

Coulomb-Tresca's condition:

$$\sigma_{\max} - \sigma_{\min} = \sigma_{o}$$

One presented it in 3D and 2D (plane stress state,  $\sigma_3 = 0$ ) as shown in Figure H4 and Figure H5



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Figure H5. Plane stress state,  $\sigma_{
m 3}$ =0, 2D



Figure H6. Isotropic material

Figure H7. Orthotropic material

Easily find the Huber-Mises condition mentioned in the law nonlinear but Coulomb-Tresca condition mentioned simple linear law in the process of analysis for the shell. If one consider the shell material with yield stress difference between tension and compression, we will show in Figure H6 & H7. To simplify further, we can assume that the orthotropic property of material considered only in tension ( $\sigma_{1}, \sigma_{2}$ > 0). Then we reshow on Figure H8. For reference, we replace ( $\varphi, \theta$ ) by (1.2). Relationship between  $\varepsilon_{i}$  and z shown in Figure H10 & H11.









Figure H10. Relationship between  $\mathcal{E}_1$  and z ESTABLISHMENT EQUATION OF YIELD SURFACE

Relationship between  $\varepsilon_{i}$ ,  $\lambda_{i}$ ,  $\chi_{i}$  and z:

$$\begin{cases} \varepsilon_1 = \lambda_1 + z\chi_1 \\ \varepsilon_2 = \lambda_2 + z\chi_2 \end{cases}$$

When the stress state is expressed as the edges of the yield surface, the legal condition is required perpendicular between ( $\varepsilon_1$ ,  $\varepsilon_2$ ) and these edges. At the corner, ( $\varepsilon_1$ ,  $\varepsilon_2$ ) will hold between two lines perpendicular to the two neighboring sides. One consider Figure H9, if the rate of deformation at a point depends on the first quadrant, one find the corresponding state of stress at this point is corner B on the yield surface. More specifically as follows

\* if 
$$\frac{-(\alpha - \beta)}{\alpha} \le \frac{\varepsilon_1}{\varepsilon_2} \le 0$$
: stress state at corner B,  
\* if  $\frac{-(\alpha - \beta)}{\alpha} \ge \frac{\varepsilon_1}{\varepsilon_2} \ge -\beta$ : stress state at corner C,  
\* if  $-\beta \ge \frac{\varepsilon_1}{\varepsilon_2} \ge -\infty$ : stress state at corner D,

One easily have

\* if 
$$\frac{h}{2} \ge z \ge -\frac{\lambda_2}{\chi_2}$$
;  $\frac{\varepsilon_1}{\varepsilon_2} > 0$ : regime B,  
\* if  $-\frac{\lambda_2}{\chi_2} \ge z \ge -\frac{\lambda_1}{\chi_1}$ ;  $\frac{\varepsilon_1}{\varepsilon_2} > 0$   
+  $0 \ge \frac{\varepsilon_2}{\varepsilon_1} \ge -\frac{(\alpha - \beta)}{\alpha}$ : regime B,

Figure H11. Relationship between  $\mathcal{E}_2$  and z

+ 
$$-\frac{(\alpha - \beta)}{\alpha} \ge \frac{\varepsilon_2}{\varepsilon_1} \ge -\beta$$
 : regime A,  
+  $-\beta \ge \frac{\varepsilon_2}{\varepsilon_1} \ge -\infty$  : regime F,  
\* if  $-\frac{\lambda_1}{\chi_1} \ge z \ge -\frac{h}{2}$ ;  $\frac{\varepsilon_1 \langle 0}{\varepsilon_2 \langle 0}$ : regime E,

Similarly one have,

\* if 
$$\frac{h}{2} \ge z \ge -\frac{\lambda_2}{\chi_2}$$
;  $\frac{\varepsilon_1}{\varepsilon_2} > 0$  : regime B,  
\* if  $-\frac{\lambda_2}{\chi_2} \ge z \ge -\frac{\lambda_1}{\chi_1}$ ;  $\frac{\varepsilon_1}{\varepsilon_2} > 0$   
+ if  $0 \ge \frac{\varepsilon_1}{\varepsilon_2} \ge -\frac{(\alpha - \beta)}{\alpha}$  : regime B,  
+ if  $-\frac{(\alpha - \beta)}{\alpha} \ge \frac{\varepsilon_1}{\varepsilon_2} \ge -\beta$  : regime C,  
+ if  $-\beta \ge \frac{\varepsilon_1}{\varepsilon_2} \ge -\infty$  : regime D,  
\* if  $-\frac{\lambda_1}{\varepsilon_2} > z \ge -\frac{h}{\varepsilon_1} \cdot \frac{\varepsilon_1}{\varepsilon_2} < 0$ 

\* if 
$$-\frac{\lambda_1}{\chi_1} \ge z \ge -\frac{n}{2}$$
;  $\frac{\varepsilon_1 \langle 0}{\varepsilon_2 \langle 0}$ : regime E,

By using this equation:  $\varepsilon_i = \lambda_i + z\chi_i$ 

$$\frac{h}{2} \ge z \ge -\left[\frac{(\alpha - \beta)\lambda_1 + \alpha\lambda_2}{(\alpha - \beta)\chi_1 + \alpha\chi_2}\right] \text{ respect to } \sigma_1 = \alpha\sigma_0$$

$$-\left[\frac{(\alpha - \beta)\lambda_1 + \alpha\lambda_2}{(\alpha - \beta)\chi_1 + \alpha\chi_2}\right] \ge z \ge -\left[\frac{(\beta)\lambda_1 + \lambda_2}{(\beta)\chi_1 + \chi_2}\right] \text{ respect to } \sigma_1 = \beta\sigma_0$$

$$-\left[\frac{(\beta)\lambda_1 + \lambda_2}{(\beta)\chi_1 + \chi_2}\right] \ge z \ge -\left[\frac{\lambda_1}{\chi_1}\right] \text{ respect to } \sigma_1 = 0$$

$$-\left[\frac{\lambda_1}{\chi_1}\right] \ge z \ge -\left[\frac{h}{2}\right] \text{ respect to } \sigma_1 = -\sigma_0$$

$$\frac{h}{2} \ge z \ge -\left[\frac{(\alpha - \beta)\lambda_1 + \alpha\lambda_2}{(\alpha - \beta)\chi_1 + \alpha\chi_2}\right] \text{ respect to } \sigma_2 = \alpha\sigma_0$$

$$-\left[\frac{(\alpha - \beta)\lambda_1 + \alpha\lambda_2}{(\alpha - \beta)\chi_1 + \alpha\chi_2}\right] \ge z \ge -\left[\frac{(\beta)\lambda_1 + \lambda_2}{(\beta)\chi_1 + \chi_2}\right] \text{ respect to } \sigma_2 = 0$$

$$-\left[\frac{(\beta)\lambda_1 + \lambda_2}{(\beta)\chi_1 + \chi_2}\right] \ge z \ge -\left[\frac{\lambda_1}{\chi_1}\right] \text{ respect to } \sigma_2 = -\sigma_0$$

$$-\left[\frac{\lambda_1}{\chi_1}\right] \ge z \ge -\left[\frac{h}{2}\right] \text{ respect to } \sigma_2 = -\sigma_0$$

In addition one have,

$$N_{1} = \int_{-\frac{h}{2}}^{\frac{h}{2}} \sigma_{1}dz; N_{2} = \int_{-\frac{h}{2}}^{\frac{h}{2}} \sigma_{2}dz; M_{1} = \int_{-\frac{h}{2}}^{\frac{h}{2}} \sigma_{1}zdz; M_{2} = \int_{-\frac{h}{2}}^{\frac{h}{2}} \sigma_{2}zdz$$

$$N_{1} = \int_{-\frac{h}{2}}^{\frac{h}{2}} -\left[\frac{(\alpha - \beta)\lambda_{1} + \alpha\lambda_{2}}{(\alpha - \beta)\chi_{1} + \alpha\chi_{2}}\right] - \left[\frac{(\beta)\lambda_{1} + \lambda_{2}}{(\beta)\chi_{1} + \chi_{2}}\right] - \left[\frac{\lambda_{1}}{\chi_{1}}\right]$$

$$N_{1} = \int_{-\frac{(\alpha - \beta)\lambda_{1} + \alpha\lambda_{2}}{(\alpha - \beta)\chi_{1} + \alpha\chi_{2}}\right] - \left[\frac{(\beta)\lambda_{1} + \lambda_{2}}{(\beta)\chi_{1} + \chi_{2}}\right] - \left[\frac{\lambda_{1}}{\chi_{1}}\right] - \left[\frac{\lambda_{1}}{\chi_{1}}\right]$$

Similarly,

$$N_{1} = \sigma_{0}h\left\{\frac{(\alpha-1)}{2} + \frac{\lambda_{1}}{h\chi_{1}} + \frac{\beta}{h}\left(\frac{\beta\lambda_{1}+\lambda_{2}}{\beta\chi_{1}+\chi_{2}}\right) + \frac{(\alpha-\beta)}{h}\left[\frac{(\alpha-\beta)\lambda_{1}+\alpha\lambda_{2}}{(\alpha-\beta)\chi_{1}+\alpha\chi_{2}}\right]\right\}$$

$$N_{2} = \sigma_{0}h\left\{\frac{(\alpha-1)}{2} + \frac{1}{h}\left(\frac{\beta\lambda_{1}+\lambda_{2}}{\beta\chi_{1}+\chi_{2}}\right) + \frac{\alpha}{h}\left[\frac{(\alpha-\beta)\lambda_{1}+\alpha\lambda_{2}}{(\alpha-\beta)\chi_{1}+\alpha\chi_{2}}\right]\right\}$$

$$M_{1} = \frac{\sigma_{0}h^{2}}{4}\left\{\frac{(\alpha+1)}{2} - 2\left(\frac{\lambda_{1}}{h\chi_{1}}\right)^{2} - \frac{2\beta}{h^{2}}\left(\frac{\beta\lambda_{1}+\lambda_{2}}{\beta\chi_{1}+\chi_{2}}\right)^{2} - \frac{2(\alpha-\beta)}{h^{2}}\left[\frac{(\alpha-\beta)\lambda_{1}+\alpha\lambda_{2}}{(\alpha-\beta)\chi_{1}+\alpha\chi_{2}}\right]^{2}\right\}$$

$$M_{2} = \frac{\sigma_{0}h^{2}}{4}\left\{\frac{(\alpha+1)}{2} - \frac{2}{h^{2}}\left(\frac{\beta\lambda_{1}+\lambda_{2}}{\beta\chi_{1}+\chi_{2}}\right)^{2} - \frac{2\alpha}{h^{2}}\left[\frac{(\alpha-\beta)\lambda_{1}+\alpha\lambda_{2}}{(\alpha-\beta)\chi_{1}+\alpha\chi_{2}}\right]^{2}\right\}$$

ILLUSTRATION

For case 
$$-\frac{h}{2} \le -\frac{\lambda_1}{\chi_1} \le -\frac{(\lambda_1 + \lambda_2)}{(\chi_1 + \chi_2)} \le -\frac{\lambda_2}{\chi_2} \le \frac{h}{2}$$
 (\*)

Presentation with dimensionless:

$$n_{1} = \frac{N_{1}}{\sigma_{0}h}; n_{2} = \frac{N_{2}}{\sigma_{0}h}; m_{1} = \frac{4M_{1}}{\sigma_{0}h^{2}}; m_{2} = \frac{4M_{2}}{\sigma_{0}h^{2}}$$

$$f = -\frac{1}{h} \left(\frac{\beta\lambda_{1} + \lambda_{2}}{\beta\chi_{1} + \chi_{2}}\right); g = -\frac{1}{h} \left[\frac{(\alpha - \beta)\lambda_{1} + \alpha\lambda_{2}}{(\alpha - \beta)\chi_{1} + \alpha\chi_{2}}\right]$$

$$p = -\frac{\lambda_{1}}{h\chi_{1}}; q = -\frac{1}{h} \left(\frac{\lambda_{1} + \lambda_{2}}{\chi_{1} + \chi_{2}}\right); r = -\frac{\lambda_{2}}{h\chi_{2}}$$

$$t = -\frac{1}{h} \left(\frac{\lambda_{1} + \beta\lambda_{2}}{\chi_{1} + \beta\chi_{2}}\right); k = -\frac{1}{h} \left[\frac{\alpha\lambda_{1} + (\alpha - \beta)\lambda_{2}}{\alpha\chi_{1} + (\alpha - \beta)\chi_{2}}\right]$$

Results are shown in Table 1 (\*). The other cases are similarly done.

Table 1 (\*).Results

Case	n <sub>i</sub>	n <sub>2</sub>	$m_i$	<i>m</i> <sub>2</sub>
$-\frac{1}{2} \le p \le q \le r \le \frac{1}{2}$ (equation (*))	$\frac{(\alpha\!-\!1)}{2}\!-\!p\!-\!(\alpha\!-\!\beta)g\!-\!\beta f$	$\frac{(\alpha-1)}{2}-f-\alpha g$	$\frac{(\alpha+1)}{2} - 2\left[p^2 + (\alpha-\beta)g^2 + \beta f^2\right]$	$\frac{(\alpha+1)}{2} - 2\left[f^2 + \alpha g^2\right]$
	$-\frac{(\alpha-1)}{2}+f-\alpha g$	$\frac{(\alpha-1)}{2} + r + (\alpha-\beta)k + \beta t$	$-\frac{(\alpha+1)}{2}+2\left[t^2+\alpha k^2\right]$	$-\frac{(\alpha+1)}{2}+2\left[r^{2}+(\alpha-\beta)k^{2}+\beta t^{2}\right]$
	$\frac{(\alpha-1)}{2} - t - \alpha k$	$-\frac{(\alpha-1)}{2}-r-(\alpha-\beta)k-\beta t$	$\frac{(\alpha+1)}{2} - 2\left[t^2 + \alpha k^2\right]$	$\frac{(\alpha+1)}{2} - 2\left[r^2 + (\alpha-\beta)k^2 + \beta t^2\right]$
	$\frac{(\alpha-1)}{2} + p + (\alpha-\beta)g + \beta f$	$\frac{(\alpha-1)}{2} + f + \alpha g$	$-\frac{(\alpha+1)}{2}+2\left[p^2+(\alpha-\beta)g^2+\beta f^2\right]$	$-\frac{(\alpha+1)}{2}+2[f^2+\alpha g^2]$
	– p + t	$\frac{(\beta-1)}{2} + r + \beta t$	$- 2 \bigl( p^2 - t^2  \bigr)$	$-\frac{(\beta+1)}{2}+2[r^2+\beta t^2]$
	$\alpha k - (\alpha - \beta)g - \beta f$	$\frac{(\beta-1)}{2} - f + (\alpha - \beta)k - \alpha g$	$2\left[\alpha k^{2}-(\alpha-\beta)g^{2}-\beta f^{2}\right]$	$\frac{(\beta+1)}{2} + 2\left[-f^2 + (\alpha-\beta)k^2 - \alpha g^2\right]$
	$-\alpha k + (\alpha - \beta)g + \beta f$	$\frac{(\beta-1)}{2}+f-(\alpha-\beta)k+\alpha g$	$\frac{(\alpha-1)}{2} - r - (\alpha - \beta)k - \beta t$	$-\frac{(\beta+1)}{2}+2[f^2-(\alpha-\beta)k^2+\alpha g^2]$
	p-t	$\frac{(\beta-1)}{2} - r - \beta t$	$2\left(p^{2}-t^{2}\right)$	$\frac{(\beta+1)}{2} - 2[r^2 + \beta t^2]$
	$\frac{(\beta-1)}{2}+t-(\alpha-\beta)g+\alpha k$	$\beta t + (\alpha - \beta)k - \alpha g$	$-\frac{(\beta+1)}{2}+2[t^2-(\alpha-\beta)g^2+\alpha k^2]$	$2\left[-\alpha g^{2}+(\alpha-\beta)k^{2}+\beta t^{2}\right]$
	$\frac{(\beta-1)}{2} - p - \beta f$	- f + r	$\frac{(\beta+1)}{2} - 2\left[p^2 + \beta f^2\right]$	$2(r^2 - f^2)$
	$\frac{(\beta-1)}{2} + p + \beta f$	t-r	$-\frac{(\beta+1)}{2}+2[p^2+\beta t^2]$	$2(-r^2 + t^2)$
	$\frac{(\beta-1)}{2} - t + (\alpha - \beta)g - \alpha k$	$-\beta t - (\alpha - \beta)k + \alpha g$	$\frac{(\beta+1)}{2} - 2\left[t^2 - (\alpha - \beta)g^2 + \alpha k^2\right]$	$2\left[\alpha g^{2}-(\alpha-\beta)k^{2}-\beta t^{2}\right]$

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## **C**ONCLUSIONS

When researching plastics for the shell we need to consider conditions for yield and yield surface of the shell in the case of orthotropic materials. This paper shows the simple way to have equation of yield surface of the shell for limit analysis method.

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