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# MHD DRAG FORCE ON WATER BASED CYLINDRICAL SHAPED ZnO NANOPARTICLE IN A CHEMICALLY REACTING NANOFLUID THROUGH CHANNEL: A THEORETICAL INVESTIGATION

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**Abstract:** An electrically conducting hydromagnetic flow of Cylindrical shaped Zinc Oxide (ZnO) and water based nanofluid is taken to investigate theoretically the behavior of various key parameter like magnetic field, heat generation, porosity, radiation and first order chemical reaction on velocity, temperature and concentration equations in a vertical channel embedded with saturated porous medium. The analytical solutions of the equations are obtained and the impact of the controlling parameters on the velocity, temperature, concentration, skin friction, Nusselt number and Sherwood number has been discussed graphically. It is concluded that with increase of  $\phi$ , both the velocity and temperature diminish. Moreover, the magnetic parameter retards the nanofluid motion whereas porosity accelerates it. Due to the presence of ZnO, the shear stresses become negative throughout the channel of porous medium. ZnO nanoparticles have used worldwide for their exclusive advances in biomedical applications (cancer cell imaging, sun screens), toxicity study, textiles and hybrid solar cells etc.

**Keywords:** ZnO-water based nanofluid, Nanoparticles volume fractions, Permeability parameter, Mass Transfer, Channel flow

## 1. INTRODUCTION

The idea of utilizing nanoparticles (alumina (Al<sub>2</sub>O<sub>3</sub>), Zinc oxide (ZnO), Copper (Cu), Silver (Ag) etc.) by adding limited quantity of nanoparticles to the base fluid (water, ethylene glycol, engine oil, grease oils etc.), which have poor thermal conductivity contrasted with that of solid nanoparticles is initiated by Choi in 1995. The pioneer work of remarkable increments in the thermal conductivity of Nanofluids, they have become the topic of lots of research. In 21<sup>st</sup> century due to the improved properties related with mass and heat transfer, anti-bacterial wetting, and spreading action, much research has been done on nanofluids. Potential uses of nanofluids incorporate transportation, cooling process like to cool automobile engines and high heat flux devices (microwave tubes etc), biomedicine, refrigeration, energy industry and so on. Now a day's ZnO is one of the most prominent metal oxide nanoparticles used in recent Nanotechnology due to its peculiar physical and chemical properties and has driven attention towards research.

Nanofluids play an important role to upgrade the heat transfer execution contrasted with pure or base fluids, so they can be considered as the most energizing heat transfer fluids. Further some detail latest exploration on nanofluid of different nanoparticles with various properties is attempted by several researchers in References (2016, 2018, 2017, 2016, 2017). Enhancement of heat transfer characteristics in the channel with trapezoidal rib–groove using nanofluids is studied by Al-Shamani et al. (2015). Hayat et al. (2015) discussed the mixed convection flow of Casson nanofluid over a stretching sheet with convectively heated chemical reaction and heat Source/Sink. Aaiza et al. (2015) observed the energy transfer in mixed convection MHD flow of nanofluid containing different shapes of nanoparticles in a channel filled with saturated porous medium.



There have been numerous investigations (2015, 2013, 2013, 2015, 2018) on the nanofluid with various effects such as heat transfer, thermal conductivity, MHD, thermal radiations, porous medium and so on. Pandey et.al (2017) studied the chemical reaction and thermal radiation effects on boundary layer flow of nanofluid over a wedge with viscous and Ohmic dissipation. MHD flow and radiation heat transfer of nanofluids in porous media with variable surface and chemical reaction is discussed by Zhang et al. (2015).Heat and mass transfer study of nanofluid with various effect is investigated by many researchers in recent years (2016, 2016). Prateek et al. (2016) reviewed the applications and properties of Zinc Oxide (Zno) nanoparticles. Recently, Dogonchi et al. (2017) studied the MHD Go-water nanofluid flow and heat transfer in a porous channel in the presence of thermal radiation effect.

This study is aimed to investigate the mixed convection MHD flow of cylindrical shape ZnO nanoparticles in water based nanofluid in a vertical permeable channel in presence of heat generation, thermal radiation and first order chemical reaction. The effects of key parameters like MHD drag force, Radiation, Chemical reaction and nanoparticle volume fraction on velocity, temperature and concentration are discussed and presented graphically as well as discussed the behaviour of skin friction, Nusselt number and Sherwood number.

#### 2. MATHEMATICAL FORMULATION

The  $\overline{\mathbf{x}}$ - axis is taken along the flow and  $\overline{\mathbf{y}}$ - axis is taken normal to the flow direction. The mixed convection is caused due to buoyancy force together with external pressure gradient applied along the  $\overline{\mathbf{x}}$ - direction. Under the usual assumption of Boussinesq approximation, the governing equations of momentum, energy and mass are as follow:

$$\rho_{\rm nf} \frac{\partial \overline{u}}{\partial \overline{t}} = -\frac{\partial \overline{p}}{\partial \overline{x}} + \mu_{\rm nf} \frac{\partial^2 \overline{u}}{\partial \overline{y}^2} + (\rho\beta)_{\rm nf} g(\overline{T} - \overline{T}_0) + (\rho\beta)_{\rm nf} g(\overline{C} - \overline{C}_0) - \left(\sigma B_0^2 + \frac{\mu_{\rm nf}}{\overline{K}_1}\right) \overline{u}, \quad (2.1)$$

$$\left(\rho C_{p}\right)_{nf} \frac{\partial \overline{T}}{\partial \overline{t}} = \kappa_{nf} \frac{\partial^{2} \overline{T}}{\partial \overline{y}^{2}} - \frac{\partial \overline{q}}{\partial \overline{y}} - Q_{0} \left(\overline{T} - \overline{T}_{0}\right) , \qquad (2.2)$$

$$\left(\frac{\partial \overline{C}}{\partial \overline{t}} + \overline{v} \frac{\partial \overline{C}}{\partial \overline{y}}\right) = D_{\rm B} \frac{\partial^2 \overline{C}}{\partial \overline{y}^2} - K_1 (\overline{C} - \overline{C}_0) \quad .$$
(2.3)

In order to encounter spherical shape of nanoparticles inside nanofluids, Hamilton and Crosser (1962) model for thermal conductivity and the Timofeeva et al. (2009) model for calculating dynamic viscosity of nanofluids have been used and they are:

$$\left\{ \begin{array}{l} n = \frac{3}{\psi}, \ \mu_{nf} = \mu_{f}(1 + a\phi + b\phi^{2}), \ (\rho)_{nf} = (1 - \phi)(\rho)_{f} + \phi(\rho)_{s}, \\ (\rho\beta)_{nf} = (1 - \phi)(\rho\beta)_{f} + \phi(\rho\beta)_{s}, \ (\rhoC_{p})_{nf} = (1 - \phi)(\rhoC_{p})_{f} + \phi(\rhoC_{p})_{s}, \\ \frac{\kappa_{nf}}{\kappa_{f}} = \frac{\kappa_{s} + (n - 1)\kappa_{f} + (n - 1)(\kappa_{s} - \kappa_{f})\phi}{\kappa_{s} + (n - 1)\kappa_{f} - (\kappa_{s} - \kappa_{f})\phi}, \\ \sigma_{nf} = \sigma_{f} \left[ 1 + \frac{3(\sigma - 1)\phi}{(\sigma + 2) - (\sigma - 1)\phi} \right] \right\}$$
(2.4)

$$\kappa_{\rm f} = \kappa_{\rm s} + (n-1)\kappa_{\rm f} - (\kappa_{\rm s} - \kappa_{\rm f})$$
  
The radiative heat flux is given by,

$$\frac{\partial \overline{q}}{\partial \overline{y}} = -4\alpha^2 \left(\overline{T} - \overline{T}_0\right)$$
(2.5)

The nondimensional parameters are:

$$\begin{cases} x = \frac{\overline{x}}{d}, y = \frac{\overline{y}}{d}, u = \frac{\overline{u}}{U_0}, t = \frac{\overline{t}U_0}{d}, v_f = \frac{\mu_f}{\rho_f}, K = \frac{\overline{K}_1}{d^2}, \lambda_n = \frac{\kappa_{nf}}{\kappa_f}, \\ \theta = \frac{\overline{T} - \overline{T}_0}{\overline{T}_w - \overline{T}_0}, C = \frac{\overline{C} - \overline{C}_0}{\overline{C}_w - \overline{C}_0}, p = \frac{\overline{p}d}{\mu U_0}, \omega = \frac{\overline{\omega}d}{U_0}, \frac{\partial p}{\partial x} = \lambda \exp(i\omega t), \\ Pe = \frac{U_0 d(\rho C_p)_f}{k_f}, Re = \frac{U_0 d}{v_f}, M^2 = \frac{\sigma d^2 B_0^2}{\mu_f}, \\ Gr = \frac{d^2 g \beta_f (T_w - T_0)}{U_0 \mu_f}, Gm = \frac{d^2 g \beta_f (\overline{C}_w - \overline{C}_0)}{U_0 \mu_f}, \end{cases}$$
(2.6)

 $\left(\begin{array}{c} R^2 = \frac{4d}{\kappa_f}, \quad Q = \frac{d}{\kappa_{nf}}, \quad S = \frac{v_0}{U_0}, \quad S = \frac{v_0}{D_B}, \quad Kr = \frac{\kappa_1 u}{U_0} \end{array}\right)$ With the help of (2.4) – (2.6), the eqs (2.1) – (2.3) reduce to the dimensionless form





$$\phi_1 \operatorname{Re} \frac{\partial u}{\partial t} = \varepsilon \lambda \exp(i\omega t) + \phi_2 \frac{\partial^2 u}{\partial y^2} - M^2 u - \frac{\phi_2}{K} u + \phi_3 \operatorname{Gr} \theta$$
(2.7)

$$\phi_4 \operatorname{Pe} \frac{1}{\lambda_n} \frac{\partial \theta}{\partial t} = \frac{\partial^2 \theta}{\partial y^2} + \frac{R^2}{\lambda_n} \theta - Q\theta$$
(2.8)

$$\frac{\partial C}{\partial t} - S \frac{\partial C}{\partial y} = \frac{1}{Sc} \frac{\partial^2 C}{\partial y^2} - KrC$$
(2.9)

The boundary conditions are

$$\begin{cases} \bar{u}(0,t) = 0, \ \bar{u}(d,t) = 0\\ \overline{T}(0,t) = T_0, \ \overline{T}(d,t) = T_{\infty}\\ \bar{C}(0,t) = C_0, \ \bar{C}(d,t) = C_{\infty} \end{cases}$$
(2.10)

The dimensionless forms of (2.10) are

$$\begin{cases} u(0,t) = 0, & u(1,t) = 0, & t > 0 \\ \theta(0,t) = 0, & \theta(1,t) = 1, & t > 0 \\ C(0,t) = 0, & C(1,t) = 1, & t > 0 \end{cases}$$
(2.11)

The eqs (2.7) and (2.9) may be re-written as

$$a_0 \frac{\partial u}{\partial t} = \epsilon \lambda \exp(i\omega t) + \phi_2 \frac{\partial^2 u}{\partial y^2} - m_0^2 u + a_1 \theta$$
(2.12)

$$b_0^2 \frac{\partial \theta}{\partial t} = \frac{\partial^2 \theta}{\partial y^2} + (b_1^2 - Q)\theta$$
(2.13)

$$\frac{\partial C}{\partial t} - S \frac{\partial C}{\partial y} = \frac{1}{Sc} \frac{\partial^2 C}{\partial y^2} - KrC$$
(2.14)

To solve the eqs  $(2.12) \sim (2.13)$  with the boundary conditions (2.11), the perturbed solutions are taken of the forms:

$$f(y,t) = f_0(y) + \varepsilon \exp(i\omega t)f_1(y)$$
(2.15)

where f stands for  $u, \theta$  or C.

On using the eqs (2.15) into eqs  $(2.12) \sim (2.14)$ , we obtain the following system of ordinary differential equations,

$$\frac{d^2 u_0(y)}{dy^2} - m_1^2 u_0(y) = -a_3 \theta_0(y) - -a_4 C_0(y), \qquad (2.16)$$

$$\frac{d^2 u_1(y)}{dy^2} - m_2^2 u_1(y) = -\lambda_1 - a_3 \theta_1(y) - a_4 C_1(y), \qquad (2.17)$$

$$\frac{d^2\theta_0(y)}{dy^2} + b_2^2\theta_0(y) = 0, \qquad (2.18)$$

$$\frac{d^2\theta_1(y)}{dy^2} + b_3^2\theta_1(y) = 0, \qquad (2.19)$$

$$\frac{d^2 C_0(y)}{dy^2} + d_0 \frac{d C_0(y)}{dy^2} - d_1 C_0(y) = 0, \qquad (2.20)$$

$$\frac{d^2 C_1(y)}{dy^2} + d_0 \frac{dC_1(y)}{dy^2} - d_2 C_1(y) = 0, \qquad (2.21)$$

The transformed boundary conditions (2.11) are

$$\begin{cases} u_0(0) = 0, \ u_0(1) = 0, \ u_1(0) = 0, \ u_1(1) = 0 \\ \theta_0(0) = 0, \ \theta_0(1) = 1, \ \theta_1(0) = 0, \ \theta_1(1) = 0 \\ C_0(0) = 0, \ C_0(1) = 1, \ C_1(0) = 0, \ C_1(1) = 0 \end{cases}$$
(2.22)

The solutions of the eqs  $(2.13) \sim (2.18)$  the under boundary conditions (2.19) yield to

$$u(y,t) = \begin{bmatrix} \left\{ B \sinh(m_1 y) + \frac{a_3 \sin(b_2 y)}{(b_2^2 + m_1^2) \sinh b_2} - \frac{a_4 A}{(n_1^2 - m_1^2)} (e^{n_1 y} - e^{-n_1 y}) \right\} \\ \left\{ \lambda_1 (\cosh(m_2 - 1)) + \lambda_1 (e^{-n_1 y}) + \lambda_1 (e^{-n_1 y}) \right\}$$
(2.23)

$$\left[ + \varepsilon \exp(i\omega t) \left\{ \frac{\lambda_1(\cosh m_2 - 1)}{m_2^2 \sinh m_2} \sinh(m_2 y) - \frac{\lambda_1}{m_2^2} (\cosh(m_2 y) - 1) \right\} \right]$$

$$\theta(\mathbf{y}) = \frac{\sin b_2 \mathbf{y}}{\sin b_2} \tag{2.24}$$

$$C(y) = A(e^{n_1 y} - e^{-n_1 y})$$
(2.25)







$$\tau = \begin{bmatrix} \frac{\partial u}{\partial y} \end{bmatrix} = \begin{bmatrix} \begin{cases} Bm_1 \cosh(m_1 y) + \frac{a_3 b_2 \cos(b_2 y)}{(b_2^2 + m_1^2) \sin b_2} \\ -\frac{a_4 An_1}{(n_1^2 - m_1^2)} (e^{n_1 y} + e^{-n_1 y}) \\ + \exp(i\omega t) \begin{cases} \frac{\lambda_1 (\cosh m_2 - 1)}{m_2 \sinh m_2} \cosh(m_2 y) \\ -\frac{\lambda_1}{m_2^2} (m_2 \cosh(m_2 y) - 1) \end{cases} \end{bmatrix}$$
(2.26)

– Nusselt Number

$$Nu = \left[\frac{\partial \theta}{\partial y}\right] = \frac{b_2 \cos b_2 y}{\sin b_2}$$
(2.27)

- Sherwood Number

$$Sh = \left[\frac{\partial C}{\partial y}\right] = An_1(e^{n_1y} + e^{-n_1y})$$
(2.28)

#### 3. VALIDITY AND ACCURACY

To check the accuracy of the present nanofluid model, comparisons of the velocity and temperature of ZnO nanofluids have been conducted with published results obtained by Aaiza et al. (2015) and presented in Tables – 1 and 2. Very good agreement between the present and the previous results confirms the accuracy of the method used. The results obtained for velocity and temperature profiles through analytically have also been weighted up in the Tables – 1 and 2, respectively, which clearly reflects the convergence of the perturbation method. It has been seen that the increasing values of  $\phi$  reduces the momentum and thermal boundary layer of ZnO nanofluids.

Table – 1: Comparison of velocity distribution with Aaiza et al. (2015) for  $\phi$  in ZnO–water based

Nationality at $N = 1, \omega = 2, N = 2, \omega = 2, N = 0.01$ .							
Aaiza et al. (2015) at $Q = 0,Gm = 0$				Present work at Q=2,Gm=2			
у	φ =0.01	φ=0.05	φ =0.09	φ=0.01	φ =0.05	φ =0.09	
0.0	0	0	0	0	0	0	
0.2	0.007568	0.006102	0.004402	0.00596	0.005589	0.00446	
0.4	0.013975	0.0111024	0.007841	0.01228	0.01101	0.00842	
0.6	0.017910	0.013637	0.009263	0.01880	0.015314	0.010843	
0.8	0.016980	0.0115310	0.007302	0.022305	0.01509	0.009567	
1.0	0	0	0	0	0	0	

Nanofluids at  $\lambda = 1, \omega = 2, R = 2, Gr = 2, K = 0.01$ :

Table – 2: Comparison of temperature distribution with Aaiza et al. (2015) for  $\phi$  in water based Nanofluids at  $\psi = 0.81$ , R = 2,  $\kappa_f = 0.613$ ,  $\kappa_s = 13$ :

Aaiza et al. (2015) at $Q = 0$				Present work at Q=5			
Y	φ =0.05	φ =0.1	<b>φ</b> =0.15	φ =0.05	φ =0.1	<b>φ</b> =0.15	
0.0	0	0	0	0	0	0	
0.2	0.36827	0.330702	0.305015	0.155128	0.14417	0.13627	
0.4	0.68811	0.624609	0.58097	0.32071	0.30108	0.28685	
0.6	0.917484	0.849022	0.80159	0.50792	0.48458	0.46750	
0.8	1.02621	0.97897	0.94584	0.72937	0.71089	0.69719	
1.0	1	1	1	1	1	1	

### 4. RESULTS AND DISCUSSIONS

The effects of flow parameters such as the Magnetic Parameter (M), Radiation parameter (R), Porosity parameter (K), Grashof numbers (Gm, Gr), Chemical reaction (Kr), Suction (S), Schmidt number (Sc) on velocity, temperature, concentration, skin-friction, Nusselt number and Sherwood number have been studied analytically and presented graphically in Fig.1 to Fig.11 for cylindrical shaped ZnO nanoparticles. Default values of the parameters which have been used in graphical analyze for the effectiveness of the model are as follows:  $\lambda = 1, t = 2, \omega = 2, R = 2, Gr = 2, Gm = 2, K = 0.01, Re = 2, Pe = 3.5, Kr = 0.05, S = 2, M = 2,$ 

 $\sigma = 0.8, \gamma = 0.62, a = 13.5, b = 904.4, Q = 2.$ 

Table – 3: Constant a and b (Empirical shape factor)					
Model	Platelet	Blade	Cylinder	Brick	
а	37.1	14.6	13.5	1.9	
b	612.6	123.3	904.4	471.4	







Figure 1: Velocity Distribution for K in ZnO-water based Nanofluid for  $\phi = 0.01$  and  $\phi = 0.1$ The nanofluids' velocity profiles are presented in Figure 1 for different estimations of the porosity parameter (K). Here velocity of the nanofluid is elevated for large values of K. In this figure  $\phi = 0.1$ is dominated by  $\phi = 0.01$  as velocity increases for small values of nanoparticles volume fraction ( $\phi$ ) and hence momentum boundary layer thickness depressed by  $\phi = 0.1$ . While momentum boundary layer thickness gradually increased by the higher porosity.

The velocity profiles (u) are plotted in Figure 2 for different estimations of nanoparticles volume fraction ( $\phi$ ) and Magnetic parameter (M). It is seen that velocity accelerates for small values of  $\phi$ . Clearly the presence of Magnetic field is to diminish the velocity in the momentum boundary layer due to the fact that the utilization of the transverse magnetic field brings about an opposing kind of force called Lorentz force, like drag force that opposes the fluid flow which results in reducing the velocity of the nanofluid in the boundary layer. Moreover, higher MHD and nanoparticle volume fraction reduced the boundary layer thick ness.



Figure 2: Velocity Distribution for  $\phi$  in ZnO-water based Nanofluid for M=0 and M=10 Figure 3 depicts the velocity profiles for different values of Grashof Number (Gr) with M. It is clearly seen that for the both values of M, an increasing in Gr elevated the nanofluid velocity. Also is noticed that M=0.5 is followed by M=5, that means M retards the naonofluid motion. The effect of heat generation parameter (Q) and radiation(R) in nanofluid velocity and temperature is presented in Figures 4 and 5. In the both Figures, with increasing Q, velocity and temperature profiles of the nanofluid decline this brings about diminishing the momentum and thermal boundary layer thickness with larger Q. It is observed that the velocity and temperature increase

with increasing radiation parameter, and hence R=2 is followed by R=0.2 in both Figures.





Significantly, without heat generation (Q = 0) maximum heat fluctuations is observed for the higher R = 2 near y = 1, while  $\theta$  becomes a linear function for the least R = 0.2.







Figure 4: Velocity Distribution for Q in ZnO-water based Nanofluid for R=0.2 and R=2



Figure 5: Temperature Distribution for Q in ZnO-water based Nanofluid for R=0.2 and R=2



Figure 6: Temperature Distribution for  $\phi$  in ZnO-water based Nanofluid for Q=1 and Q=5





Figure 6 is a graphical representation of different values of nanoparticles volume fraction ( $\phi$ ) and heat generation (Q) on temperature distribution. It is seen that temperature of ZnO diminishes for large values of  $\phi$  and Q and hence the thermal boundary layer thickness becomes smaller. Moreover, the greater heat generation always suppresses the thermal boundary layer thickness. However, maximum  $\theta$  is occurred near y = 1 for the least Q = 1 and this behaviour is opposite for the higher Q = 5.



Figure 7: Concentration Distribution for Kr in S=0.2 and S=2

Figure 7 demonstrates the behavior of chemical reaction (Kr) and suction (S) on concentration profiles. It is marked that the concentration boundary layer is diminished by the effect of maximum suction parameter (S=2) and destructive chemical reaction (Kr = 1.5). The behavior of Kr and S reduces the species concentration in the channel of porous medium and hence S=2 is dominated by S=0.2 in a chemically reacting fluid.



Fig.8: Concentration Distribution for Sc in air at 25° Celsius

The behavior of Sc on concentration profiles for air at the temperature of  $25^{\circ}$  Celsius is presented in Figure 8. An increase of Sc from Ammonia (Sc=0.57), Oxygen (Sc=0.84) through Methane (Sc=0.99) to Carbon dioxide (Sc=1.14) depress the species concentration boundary layer as heavy molecular diffusivity suppress the concentration profiles in boundary layer.



Figure 9: Skin friction for Gm in ZnO-water based Nanofluid for M=0.5 and M=5





In Figure 9, Skin friction profile for mass Grashof number (Gm) in ZnO-water based Nanofluid is plotted for two different M=0.5 and M=5 against nanoparticle volume fraction ( $\phi$ ). Skin friction profiles gradually diminishes, when Gm and  $\phi$  increase. In presence of ZnO, the shear stresses become negative at the plate y = 0 for higher values of Gm=4, 6 and M=5. At  $\phi$  =0.05 and Gm=2 a reverse behavior is observed for the shear stresses.



Figure 10: Rate of heat transfer at y=0 for Q and R in ZnO-water based Nanofluid when  $\phi$ =0.01 and  $\phi$ = 0.5 The effects of heat generation and radiation parameter on Nusselt number is displayed in Figure 10. Due to the bigger thermal conductivity of ZnO, the rate of heat transfer has dominant value at the lower  $\phi$  = 0.01 than the higher value  $\phi$  = 0.5. However, the thermal radiation (R) enhanced the rate of heat transfer, but this behaviour is reversed for the heat generation (Q).



Figure 11: Sherwood number for Sc and Kr

In Figure 11 the effect of endothermic first order chemical reaction (Kr) and different values of Schmidt number (Sc) against suction (S) is presented for Sherwood number profile. Due to the cause of endothermic rate, the concentration is reduced throughout the channel and it is observed that Kr=2 is dominated by Kr=0.2. Sherwood number profiles strongly reduce with increasing Sc values. Also Sherwood number diminishes for greater values of suction parameter.

## 5. CONCLUSIONS

Impact of nanoparticle volume fraction of water based ZnO nanofluid in a chemically reacting fluid through a channel of porous medium with magnetic drag force and thermal radiation is investigated. During investigation of the results obtained through graphs and tables, few significant conclusions are mentioned below:

- The results obtained here are useful in engineering problems, particularly, in energy systems, rheology, material processing, lubrication and biomedical applications
- Due to Cylindrical shaped water based Zinc Oxide nanofluid, the velocity overshoot is observed in all the cases.
- The behaviour of nanoparticle volume fraction  $\phi = 0.01$  is the most dominating agent in Zinc Oxide nanofluid over  $\phi = 0.1$  for velocity, temperature and rate of heat transfer.
- In presence of water based ZnO nanofluid, porosity and natural convection enhanced the boundary layer thickness in the channel.
- —Shear stresses at y=0 gradually become negative in presence of Gm,  $\phi$  and M.





Significantly, the rate of mass transfer at y = 0 is more dominant at Kr = 0.2 than Kr = 2.0 and reduces exponentially to zero with the effects of Sc and S;

- -Higher radiation R = 2 plays a dominant agent over the lower radiation R = 0.2 on velocity and temperature of ZnO nanofluid, while they are reduced by higher Q.
- Lower magnetic field (M = 0) and heat generation (Q = 1) boost the profiles of velocity and temperature of ZnO nanofluid in comparison of higher M and Q, while  $\phi$  reduces the thickness of boundary layer of u and  $\theta$ .

Nomenclature:	

ū	dimensional velocity	κ <sub>nf</sub>	thermal conductivity of	φ	solid volume fraction of
	component		nanofluid		nanoparticles
Т	dimensional temperature of the	$\beta_{nf}$	coefficient of thermal	K <sub>1</sub>	dimensional permeability
	nanofluid		expansion of nanofluid		parameter
С	dimensional concentration of	$\rho_{nf}$	density of the nanofluid	$Q_0$	dimensional additional heat
	the nanofluid				source
u	non-dimensional velocity	$\mu_{nf}$	viscosity of the nanofluid	Q	non- dimensional additional
	component along x-axis				heat source
θ	non-dimensional temperature	$\sigma_{\rm nf}$	electric conductivity of	Pe	Peclet Number
	of the nanofluid		nanofluid		
ψ	non-dimensional concentration	$(\rho C_p)_{n}$	heat capacitance of the	К	non-dimensional permeability
	of the nanofluid		nanofluid		parameter
$\overline{T}_0$	temperature at $y = 0$	$(\rho\beta)_{nf}$	thermal expansion	Re	Renolds Number
Ű			coefficient of the nanofluid		
Ē	concentration at $y = 0$	<i>a</i> .	thermal diffusivity of the	c	substitute(S > 0)
$C_0$	concentration at $y = 0$	$a_{nf}$	nanofluid	3	successfully (5 > 0),
_					injection $(5 < 0)$
Tw	temperature of the wall	$\rho_{f}$	density of the base-fluid	М	magnetic parameter
$\bar{C}_{w}$	concentration of the wall	$\rho_s$	density of the nanoparticles	Kr	chemical reaction parameter
B <sub>0</sub>	magnetic field strength	$\mu_{f}$	viscosity of the base-fluid	Sc	Schmidt number
g	acceleration due to gravity	$\rho_{f}$	density of the base-fluid	Gr	thermal Grashof number
κ <sub>f</sub>	thermal conductivity of base-	ρ <sub>s</sub>	density of the nanoparticles	Gm	mass Grashof number
	fluid				
κ <sub>s</sub>	thermal conductivity of	$(\rho\beta)_{f}$	thermal expansion	3	reference constant due to
	nanoparticles		coefficient of the base-fluid		perturbation
$\sigma_{\rm f}$	electric conductivity of base-	$(\rho\beta)_{s}$	thermal expansion	ω	angular velocity
	fluid	_	coefficient of the		
			nanoparticles		
t	time	$(\rho C_p)_f$	heat capacitance of the	a,	empirical Shape factor
			base-fluid	b	
γ	Sphericity	$(\rho C_n)_s$	heat capacitance of the		
l '	-	(1 P)3	nanoparticles		

Appendix:

$$\begin{cases} \varphi_{1} = (1 - \varphi) + \varphi \frac{\rho_{s}}{\rho_{f}}, \ \varphi_{2} = 1 + a\varphi + b\varphi^{2}, \varphi_{3} = 1 + \frac{3(\sigma - 1)\varphi}{(\sigma + 2) - (\sigma - 1)\varphi} \\ \varphi_{4} = (1 - \varphi) + \varphi \frac{(\rho\beta)_{s}}{(\rho\beta)_{f}}, \ \varphi_{5} = (1 - \varphi) + \varphi \frac{(\rhoC_{p})_{s}}{(\rhoC_{p})_{f}}, \ \lambda_{1} = \frac{\lambda}{\varphi_{2}}, \\ a_{0} = \varphi_{1}Re, \ m_{0}^{2} = M^{2} + \frac{\varphi_{2}}{K}, \ a_{1} = \varphi_{4}Gr, \ a_{2} = \varphi_{4}Gm, \ a_{3} = \frac{a_{1}}{\varphi_{2}}, a_{4} = \frac{a_{2}}{\varphi_{2}} \\ b_{0}^{2} = \varphi_{5}Pe\frac{1}{\lambda_{n}}, \ b_{1}^{2} = \frac{R^{2}}{\lambda_{n}}, b_{2}^{2} = b_{1}^{2} - Q, b_{3}^{2} = b_{2}^{2} - b_{0}^{2}i\omega, \ m_{1}^{2} = \frac{m_{0}^{2}}{\varphi_{2}}, \ m_{2}^{2} = \frac{m_{0}^{2} + a_{0}i\omega}{\varphi_{2}}, \\ d_{0} = Sc.S, \ d_{1} = Sc.Kr, d_{2} = Sc(i\omega + Kr), \ n_{1} = \frac{d_{0} + \sqrt{d_{0}^{2} + 4d_{1}}}{2}, A = \frac{1}{e^{n_{1}} - e^{-n_{1}}}, \\ B = \frac{1}{\sinh m_{1}} \left\{ \frac{a_{4}A(e^{n_{1}} - e^{-n_{1}})}{(n_{1}^{2} - m_{1}^{2})} - \frac{a_{3}}{b_{2}^{2} + m_{1}^{2}} \right\}$$

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